The Pennsylvania State University

The Graduate School

Department of Geosciences

INVESTIGATING THE PALEOCENE-EOCENE CARBON-CYCLE PERTURBATION:

AN EARTH SYSTEM MODEL APPROACH

A Thesis in

Geosciences

by

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Submitted in Partial Fulfillment of the Requirements for the Degree of

Doctor of Philosophy

August 2007

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ABSTRACT

The Paleocene-Eocene Thermal Maximum, or PETM, was a time of sudden and major environmental change 55 million years ago, linked to a ~200 ka interval of global warming. It has been difficult to identify the cause of the PETM, however, because the <10 ka onset of the event is too rapid to explain by processes that operate on geologic timescales. This project was directed at using changes in the marine record of the PETM, a negative δ^{13} C excursion, and a decrease in CaCO₃ wt%, along with an Earth system model to better understand the cause and characteristics of the PETM.

A two-fold approach is presented here that utilizes the capabilities of the GENIE-1 Earth system model to simulate the spatial distribution of $CaCO_3$ on the sea floor, and to generate synthetic sediment cores. Chapter 2 presents a comparison of simulated patterns of sea floor $CaCO_3$ distribution with the observed PETM distribution of $CaCO_3$. Pulses of CO_2 ranging from 2,300 Pg C to 70,000 Pg C were added to the model Late Paleocene Earth system to simulate the isotopic end-members of potential sources, biogenic methane ($\delta^{13}C$ of -60‰) and mantle CO_2 ($\delta^{13}C$ of -5‰). A pulse of 7,000 Pg C came closest to reproducing observations when variable rates of sediment mixing by bioturbation were permitted.

In Chapter 3 the δ^{13} C and sediment records generated from these simulations were analysed to determine what processes affect the final δ^{13} C excursion that is recorded in sediments, and how well this excursion refects the δ^{13} C of surface ocean DIC. Key

factors determining the magitude of the $\delta^{13}C$ excursion recorded in sediments are the rate of recovery of the saturation state of seawater, and sediment mixing by bioturbation.

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ACKNOWLEDGEMENTS

I knew I was in trouble when they stopped using numbers.
-Anonymous mathphobe

Sincere thanks go to:

Lee Kump for invaluable advice and insight, patience, and for disagreeing frequently and energetically;

Klaus Keller for mentorship, friendship, and calling it as he sees it;

Andy Ridgwell for giving of his time and expertise with seemingly infinite generosity (though he should really take more vacations);

Tim Bralower and Ray Najjar for a thorough reading of the manuscript and detailed annotations

Matt, Katja, Burt, Chris, Lisa, Geoff, Tyrone, and Scott for friendship, thoughtful conversations, and pleasant coffee breaks;

Pearl, Irene, and Niels for taking good care of me;

and Jeff for gleefully finding my spelling mistakes, and patiently waiting for me to not be a student anymore.

Chapter 1

Overview

The Paleocene-Eocene boundary interval was a time of sudden and major environmental change 55 million years ago, linked to a short interval (~200 ka, Rohl et al., 2000) of global warming, referred to as the Paleocene-Eocene Thermal Maximum, or PETM. Evidence of PETM warming in the rock record is accompanied by carbon isotope excursions (CIEs) of -2.5 to -4‰ in the marine realm (e.g., Kennett and Stott, 1991; Cramer et al., 1999; Zachos et al., 2003; Fig. 1-1), and -5 to -7‰ in the terrestrial realm (e.g., Bowen et al., 2004 and references therein), suggesting that the warming was caused by a pulse of carbon input with a low carbon isotope composition (δ^{13} C). This is consistent with another feature of the PETM, widespread dissolution of CaCO₃ in marine sediments (e.g., Zachos et al., 2005; Murphy et al., 2006; Fig. 1-1).

The PETM has generated much interest for two main reasons: First, it is viewed as an analogue of the kind of climate change event that is being driven by anthropogenic emissions of greenhouse gasses. Second, it has defied explanation by the usual means. Often, negative $\delta^{13}C$ excursions can be explained by the transfer of carbon from organic to inorganic sedimentary carbon reservoirs (e.g., an increase in weathering of fossil organic matter on land and its ultimate remineralization to dissolved inorganic carbon in seawater). At the PETM, however, the onset of the $\delta^{13}C$ excursion was rapid, perhaps even less than 10 ka (e.g., Kennett and Stott, 1991), an unreasonably short period of time to transfer the large amount of carbon required to generate such a large and global $\delta^{13}C$

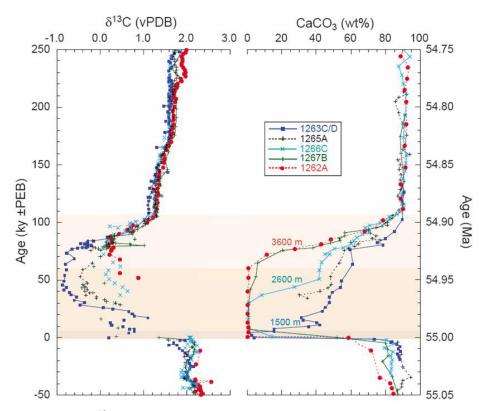


Figure 1-1: PETM δ^{13} C and CaCO₃ wt% records from Walvis Ridge in the southern Atlantic Ocean. Dark, medium, and light shading represent the onset, body, and recovery of the event, respectively. From Zachos et al. (2005).

perturbation through slow geological mechanisms. Dickens et al. (1995) framed the problem thus: to generate a -3‰ δ^{13} C excursion in the modern ocean and atmospheric carbon reservoir over 10 ka, mass balance calculations indicate that more organic carbon (δ^{13} C ~-25‰) would have to be buried in the ocean than presently exists in the terrestrial carbon reservoir. Furthermore, if the excursion were to be generated by volcanic degassing of mantle CO_2 (δ^{13} C ~-5‰), degassing would have had to occur at a rate 125 times greater than estimated long-term rates. The only reasonable alternative, these authors suggested, was that the carbon came from biogenic methane (δ^{13} C ~-60‰) and

that the methane was derived from thermal dissociation of methane clathrates within marine sediments. Other sources of carbon have been suggested since then, including organic carbon, both terrestrial from the burning of massive Upper Paleocene peat deposits (Kurtz et al., 2003), and marine, from desiccated epeiric sea sediments (Higgins and Schrag, 2006); thermogenic methane generated by sill emplacement within organic-rich sediments during North Atlantic Igneous Province volcanism (Svensen et al., 2004; Storey et al., 2007); and even extraterrestrial carbon from cometary impact (Kent et al., 2003).

Mass balance arguments can be used to determine how much carbon from a particular source is required to generate the CIE, but mass balance cannot provide a unique solution. There are many possible combinations of $\delta^{13}C$ and carbon flux magnitudes that could generate the observed $\delta^{13}C$ excursion, especially if multiple sources are considered. Further constraints are required to narrow down the options. This project was directed at using the observed CaCO₃ wt% distribution from below and within the PETM interval to guide an Earth system model configured for the Late Paleocene.

The tool used for this study is the Earth system model GENIE-1 (Grid ENabled Integrated Earth system model). The version of GENIE-1 used here (Lenton et al., 2006) comprises the coupled ocean-atmosphere model C-GOLDSTEIN (Edwards and Marsh, 2005 and references therein), a model of marine biogeochemical cycling (Ridgwell, 2001; Ridgwell et al., 2007), and a sediment model that can predict the CaCO₃ content of sea floor (coretop) sediments, has the option to simulate bioturbation, and can generate a

synthetic sediment column for direct comparison with marine cores (Ridgwell and Hargreaves, 2007).

There were several advantages to GENIE-1 over other models that led to its use for this project. GENIE-1 is a fully 3-D ocean model, but because of simplified frictional geostrophic physics and a relatively coarse resolution it runs orders of magnitude faster than other models initially considered for this study, including the Bern 2.5D model and an EMIC to be built in-house with the Geophysical Fluid Dynamic Lab's Modular Ocean Model 3 (MOM3) at its core. GENIE-1 also has great flexibility with respect to computing platforms. A further advantage of GENIE-1 was the coincidence of objectives between this study, and the developers of the coupled ocean-atmosphere component of GENIE-1, C-GOLDSTEIN, at the National Oceanographic Center in Southampton, UK. These researchers were interested in pursuing simulations with paleogeographic reconstructions, and were looking specifically at the ocean circulation at the Paleocene-Eocene boundary. As a result of using GENIE-1, this project also benefited enormously from the assistance of Andrew Ridgwell, who was developing the marine biogeochemical and sediment chemistry portions of the model along the lines of his Ph.D. thesis model, SUE (Ridgwell, 2001) at the beginning of our collaboration. Andrew Ridgwell worked on upgrades to the model throughout the course of the study, including development of synthetic model sediment cores. This capability allows a far more direct comparison between model results and data, it is certain to provide valuable insights into how diagenetic processes are recorded in sea floor sediment cores.

The outcome of this project is two manuscripts. The first manuscript presents a comparison of observed PETM CaCO₃ wt% with model core top CaCO₃ wt% values for

eight carbon release scenarios. Carbon release scenarios were selected to span the range of δ^{13} C of possible carbon sources. The low δ^{13} C end-member is biogenic methane (δ^{13} C of -60%), and the high δ^{13} C end-member is mantle CO₂ (δ^{13} C of -5%). Intermediate scenarios include δ^{13} C of -22‰, representative of organic carbon from peat (Kurtz et al., 2003), and δ^{13} C of -12, -9.5, and -6‰, reflecting combinations of high δ^{13} C carbon derived from mantle CO_2 and low $\delta^{13}C$ carbon as from biogenic methane, organic carbon, or thermogenic methane (δ^{13} C of -35 to -50%, Hunt, 1996). Although the average PETM CIE has been cited as -2.5 to -3% (e.g., Dickens et al., 1995; Pagani et al., 2006), a survey of published δ^{13} C analyses conducted for this project suggests that an excursion of -4‰ is more appropriate. The requisite carbon pulses for carbon sources with $\delta^{13}C$ of -60, -22, -12, -9.5, -6, and -5\% are 2,300, 7,000, 15,000, 21,000, 46,000, and 70,000 Pg C, respectively. Each scenario was run with bioturbation, and the 2,300 and 7,000 Pg C scenarios were also run without bioturbation. Bioturbation is an important consideration because it would have re-supplied the surface sediments with CaCO₃ even as the CO₂ from the PETM pulse was dissolving CaCO₃ in those sediments. Thus, sediments that were bioturbated would have maintained higher wt% CaCO₃ than non-bioturbated sediments. Of the pulse sizes modeled in this study, 7,000 Pg C is the smallest that can reasonably reproduce the observed distribution of CaCO₃ during the PETM, but only if we allow that shallow sediments in the central equatorial Pacific were bioturbated during the PETM, and that deeper sediments in the Pacific and all sediments at Walvis Ridge in the southern Atlantic Ocean were not.

The second manuscript addresses a more basic issue: how well does the sedimentary record represent the actual δ^{13} C history of dissolved inorganic carbon (DIC) in seawater? This manuscript was motivated in part by the discrepancy between terrestrial and marine records of the PETM CIE. Which best reflects the excursion that actually occurred in the atmosphere: the -5 to -7% excursion on land, or the -4% excursion in the sea? Perhaps the marine sediments recorded smaller excursions because of the dissolution that occurred at the PETM. In this case, variable dissolution might explain the range of marine δ^{13} C excursions: at some sites the CIE is -4% or larger, whereas at others it is little more than -1%. This would lead us to expect that the largest CIE magnitudes would occur at sites with the least dissolution. However, the data compiled for this study do not show a simple relationship between excursion magnitude and the change in wt% CaCO₃. In general, more of the larger excursions occur in sediments with a relatively small decrease in wt% CaCO₃, but some sites with near complete dissolution of CaCO₃ still record excursions of similar magnitude to sites with hardly any dissolution. Insights from modeling are used along with model sediment cores and δ^{13} C and CaCO₃ wt% observations to further elaborate on the nature of the PETM event.

The first of the two manuscripts was co-authored with Andy Ridgwell and Lee Kump and is being revised for further consideration by *Geology*. The second manuscript will be submitted to *Paleoceanography* and is co-authored with Andy Ridgwell, Lee Kump, and Tim Bralower. Andy provided assistance with modelling, and toward the end of the project took over running the final simulations with a new model version. Andy, Lee, and Tim have provided editorial assistance, and Tim has provided valuable guidance

in sorting through the literature on the PETM record in marine sediments. My contribution included taking the lead in writing both manuscripts, running the majority of the simulations, performing all post-processing tasks, and compiling PETM observations from the literature.

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Chapter 2

The sedimentary response to Paleocene Eocene Thermal Maximum carbon release: A model-data comparison¹

Abstract

An Earth system model of intermediate complexity configured with early Eocene paleogeography is used to constrain possible sources of carbon that may have caused global warming at the Paleocene-Eocene boundary. We find that 7,000 Pg C with δ^{13} C of -22‰ is the smallest pulse modeled here to reasonably reproduce observations of seafloor CaCO₃ dissolution. This pulse is unlikely to have been the result of methane hydrate dissolution or mantle CO₂ from volcanic degassing alone, suggesting that additional sources of CO₂ such as from the oxidation of sedimentary organic carbon or thermogenic methane must also have contributed. Observed contrasts in dissolution intensity between Atlantic and Pacific sites are reproduced in the model by reducing bioturbation in the Atlantic, presumably associated with the spread of low-oxygen bottom waters.

¹ K. Panchuk, A. Ridgwell, and L. Kump, for submission to Geology.

Introduction

The Paleocene-Eocene boundary (55 Ma) was characterized by geologically abrupt global warming, with ocean temperatures rising by ~5 °C over as little as 10 ka (e.g., Kennett and Stott, 1991; Zachos et al., 2003; Sluijs et al., 2006). Marine sediments deposited during the Paleocene-Eocene Thermal Maximum (PETM) recorded an average carbon isotope (δ^{13} C) excursion of -3% (e.g., Kennett and Stott, 1991; Zachos et al., 2003), as well as a marked decrease in calcium carbonate weight percent (CaCO₃) wt%; e.g., Thomas, 1998; Thomas et al., 1999; Zachos et al., 2005). The actual excursion was likely closer to -4% (e.g., Pagani et al., 2006), as recorded in well-preserved planktonic and benthic forams (Lu and Keller, 1993; Cramer et al., 1999; Thomas et al., 2002), although -5 to -7% excursions in terrestrial records suggest that it might have been even higher (e.g., Bowen et al., 2004 and references therein). In contrast, smaller (-2.5%) excursions have been measured in bulk carbonates at Walvis Ridge (Zachos et al., 2005), where sediments undergo near complete dissolution. Together, the negative δ¹³C excursion and CaCO₃ dissolution are consistent with the PETM being caused by a large pulse of ¹³C-depleted carbon to the ocean or atmosphere (Dickens et al., 1995; Thomas and Shackleton, 1996). For this reason, the PETM has been viewed as an important possible analogue for the future consequences of fossil fuel burning.

The quantity of C needed from a potential source is constrained by its characteristic isotopic signature and the magnitude of the isotopic excursion, but these constraints cannot by themselves identify the actual source. Fortunately, more information is provided by the observed climatic and sedimentary response during the

PETM. Our approach to identifying the carbon source responsible for the PETM is to compare observed spatial and temporal variations in the CaCO₃ content of marine PETM sediments with predicted changes in the spatial distribution of CaCO₃ for three potential PETM scenarios: 1) biogenic methane (i.e., the methane clathrate hypothesis of Dickens et al., 1995), with δ^{13} C of -60%; 2) rapid oxidation of sedimentary organic carbon (including peat as in Kurtz et al., 2003), with δ^{13} C of -22%; and 3) massive addition of mantle-derived volcanic CO₂ with δ^{13} C of -5%. These three scenarios do not represent all possible sources or combinations of sources, but they do span the range of source δ^{13} C values that have been suggested in the literature (e.g., Kent et al., 2003; Svensen et al., 2004; Storey et al., 2007).

This approach has only recently become possible with the development of intermediate complexity Earth system models. We use the GENIE-1 model (Lenton et al., 2006; Ridgwell et al., 2007; www.genie.ac.uk) to simulate climate, atmospheric composition, and marine biogeochemistry within a paleogeographically appropriate configuration while accounting for basic climate feedbacks. GENIE-1 is particularly valuable because the predicted spatial distribution of CaCO₃ burial and sea floor CaCO₃ δ^{13} C can be compared directly to geologic data describing the state of the Earth system before and during the PETM.

Methods

GENIE-1 includes the fast climate model of Edwards and Marsh (2005), consisting of a reduced physics (frictional geostrophic) 3-D ocean circulation model

(Edwards, 1996), coupled to a 2-D energy-moisture balance model (EMBM) of the atmosphere and a dynamic-thermodynamic sea-ice model based on that of Weaver et al. (2001). The ocean model includes a representation of marine carbon cycling (BIOGEM; Ridgwell et al., 2007), and, particularly important for interpreting the marine geological record, there is also a model of sea floor biogenic carbonate preservation (SEDGEM; Ridgwell and Hargreaves, 2007). In the model the CaCO₃ export from the ocean surface layer depends upon the CaCO₃:POC rain ratio parameter, the amount of CaCO₃ relative to particulate organic carbon (POC) exported as biogenic detritus from the surface ocean. POC export is itself a function of the PO₄³- concentration in the ocean surface layer using a Monod growth formulation. POC and CaCO₃ are remineralized in the water column according to exponential decay functions. All POC reaching the sea floor is remineralized and returned to the ocean, however CaCO₃ which reaches the ocean floor is passed to the surface sediment layer of the sediment model. SEDGEM uses a look-up table generated by the model of Archer (1991) to determine how much sedimentary CaCO₃ dissolves. CO₂ evolved from the remineralization of POC is available to dissolve CaCO₃ within the sediments.

Sediment layers within SEDGEM are 1 cm thick. If the top layer is filled, SEDGEM moves the sediments down one layer. If all CaCO₃ in the top layer is dissolved, it is replaced by moving sediment layers up the column. In this way SEDGEM synthesizes a stratigraphic record that can be affected by the long-term depositional history of sea floor sediments. GENIE-1 was spun up for 150 ka to allow the system to reach geochemical equilibrium between weathering and sedimentation and to build up a sufficiently deep sedimentary column.

SEDGEM also allows for sediments to be bioturbated. Bioturbation is represented by mixing adjacent layers within the sediment stack. The rate of mixing is highest at the top of the stack, at 16 cm²/ka, and decreases down the stack of sediment layers according to an e-folding depth of 1 cm (Peng et al., 1979; Ridgwell, 2001). Bioturbation is important to consider because it acts as a source of CaCO₃ in the sediment surface layer, increasing the buffering capacity of the sediments for a given size of carbon pulse. Without bioturbation, the amount of CaCO₃ available for dissolution is limited by the initial CaCO₃ content of the sediments and the effective depth of diffusion. As CaCO₃ dissolves, detrital (non-carbonate) material accumulates and eventually forms a cap which prevents further dissolution of CaCO₃ deeper within the sediment column. The detrital flux to the sediments is 0.18 g cm⁻²/ka, a value tuned to better approximate Latest Paleocene CaCO₃ wt% distributions. The non-carbonate detrital flux used here is still very similar to the value typical of the modern Pacific Ocean that was used in previous long-term fossil-fuel CO₂ removal simulations by Archer et al. (1998), 0.16 g cm⁻²/ka.

The model has the early Eocene bathymetry and continental configuration of Bice et al. (1998; Appendix B) and the zonally averaged wind stress field of Bice and Marotzke (2002), both transformed to a 36×36 equal area model domain. A Late Paleocene pCO₂ level of 750 ppm was chosen because it gives model global average deep ocean temperatures close to temperatures derived from δ¹⁸O analyses of benthic foraminifera, typically 11-12°C (Kennett and Stott, 1991; Bralower et al., 1995; Thomas et al., 2002), after the climate and ocean carbon cycle are equilibrated for 10,000 model years. An atmospheric CO₂ concentration of 750 ppm is within the range of proxy

studies and numerical models which have yielded estimates of Late Paleocene pCO₂ levels ranging from 300 ppm to >2000 ppm (Shellito et al., 2003 and references therein; Pagani et al., 2006).

The global weathering rate of HCO₃⁻ to the ocean and the CaCO₃:POC rain ratio are important sources of uncertainty because these parameters control the model CaCO₃ wt% distribution. We constrain these uncertainties by comparing the CaCO₃ wt% distribution from an ensemble of simulations reflecting a range of weathering rates and spatiotemporally uniform rain ratios to the observed Late Paleocene CaCO₃ wt%. All observations and their sources are summarized in Appendix B. It should be noted, however, that available data are biased toward carbonate-rich sediments because those sediments are often specifically targeted for drilling.

Another consideration when comparing data to model results is that CaCO₃ wt% is topographically controlled, so the coarse bathymetric grid means that data sites do not always correspond to model cells of the appropriate depth. There is also substantial uncertainty associated with sediment paleodepths and with the reconstructed location of Late Paleocene spreading ridges. Thus, we would not *a priori* expect to achieve a perfect model-data fit at all sites. We address this problem for both initial conditions and PETM simulations by plotting CaCO₃ wt% versus depth for a subset of grid cells surrounding data-rich sites. The best match is determined by finding the ensemble member with the minimum root mean square error (RMSE). This process is detailed in Appendix C.

We assess the geochemical impacts of three possible carbon sources: the "methane" scenario, in which biogenic methane from clathrates (-60‰) is presumed to oxidize immediately to CO₂; the "organic C" scenario (-22‰); and the "volcanic CO₂"

scenario (-5‰). The amount of CO_2 added in each case is that which instantaneously decreases the $\delta^{13}C$ of the combined atmosphere and ocean carbon reservoir by 4‰. Following Dickens (2003), CO_2 was applied uniformly to the atmosphere over 10 ka, then the model was run for an additional 100 ka without CO_2 flux forcing. For the methane, organic C, and volcanic CO_2 scenarios, 2,300, 7,000, and 70,000 Pg C, respectively, were added. The methane and organic C scenarios were run twice: once without bioturbation of the sediments, and once with bioturbation. Three intermediate scenarios with bioturbation on were also run to populate the gap between 7,000 Pg C and 70,000 Pg C. Pulses of 15,000, 21,000, and 46,000 Pg C were added, reflecting carbon sources with $\delta^{13}C$ of -12‰, -9.5‰, and -6‰, respectively.

Results and Discussion

Initial Late Paleocene conditions

Observations suggest that the pre-PETM calcite compensation depth (CCD), the depth at which CaCO₃ is no longer found in sediments, was ~3500 to 4000 m or below in the data-rich regions of the central equatorial Pacific, Walvis Ridge in the Atlantic Ocean, and the southern Indian Ocean (regions a, b, and c, respectively in Fig. 2-1). The weathering flux – rain ratio ensemble member best fitting the observations has a CaCO₃:POC rain ratio of 0.175, somewhat higher than estimates of the modern global average (e.g., 0.06, Sarmiento et al., 2002; 0.14, Ridgwell et al., 2007), and a total global weathering rate of 50 Tmol HCO₃ a⁻¹, which is larger than the 23-39 Tmol HCO₃ a⁻¹

modern estimate (Munhoven, 2002) but less than Early Eocene model simulations (>70 Tmol HCO₃⁻ a⁻¹, Gibbs et al, 1999). Some specific sites, particularly in the Indian Ocean, are difficult to reproduce. This may reflect regional deficiencies in ocean circulation and/or biogeochemistry, or the fact that the coarse bathymetry does not permit sufficiently detailed representations of rapidly sloping regions.

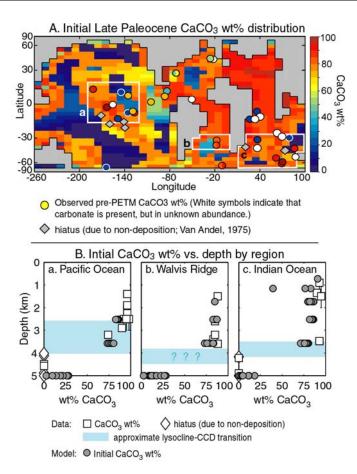


Figure 2-1: Pre-PETM model CaCO₃ wt% plotted in map view (A) and versus depth (B), both with observations superimposed. Results in B are plotted for data-rich regions delineated in A by white boxes. These are: a) the central equatorial Pacific Ocean, b) Walvis Ridge, and c) the southern Indian Ocean. See Appendix A for data values and sources.

The circulation state of the deep ocean has been traced through gradients of benthic foraminiferal δ^{13} C (e.g., Nunes and Norris, 2005; Tripati and Elderfield, 2005). With atmospheric CO_2 δ^{13} C set at -6% we obtain model bottom water δ^{13} C values which are consistent with those compiled by Nunes and Norris (2006; Fig. 2-2). Data and our model results show a decreasing δ^{13} C gradient from the southern Atlantic toward the northern Atlantic, and from the Indian Ocean toward the Pacific, from which a dominant Southern Ocean deep-water source region has been inferred for the Late Paleocene. Nd isotopic records also support this interpretation, and suggest an intermediate water source in the North Pacific (Thomas et al., 2003) which is predicted by the model as well. Thus, despite the low spatial resolution of the ocean model and lack of a dynamic atmosphere, initial model circulation patterns are not inconsistent with data-based interpretations of Late Paleocene circulation patterns.

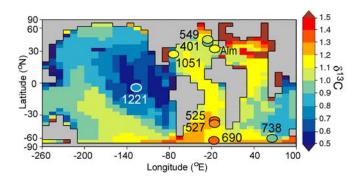


Figure 2-2: Model bottom water DIC δ^{13} C with observations as compiled by Nunes and Norris (2006) superimposed (circular markers).

PETM results

Globally averaged whole ocean temperatures increased by ~1.5, 4, and 14°C for the methane, organic C, and volcanic CO₂ cases, respectively (Fig. 2-3B), corresponding to increases in atmospheric CO₂ from 750 to 1,140, 2,200, and 25,500 ppm, respectively (Fig 2-3C). The larger the CO₂ pulse, the more sea floor CaCO₃ wt% decreases (Fig.2-3D). In the case of the 70,000 Pg C volcanic CO₂ pulse, CaCO₃ is completely dissolved over the entire sea floor, severely over-predicting CaCO₃ dissolution. The effect of bioturbation is apparent for the 2,300 and 7,000 Pg C pulses. When sediments are bioturbated (dashed lines in Fig. 2-3D) globally averaged CaCO₃ wt% is higher at the PETM than for non-bioturbated sediments (solid lines in Fig. 2-3D), and the peak PETM atmospheric CO₂ level is slightly lower for bioturbated sediments than non-bioturbated sediments (black lines in Fig. 2-3C), reflecting the enhanced buffering capacity afforded by a continuous re-supply of CaCO₃ to the surface sediment layer by bioturbation.

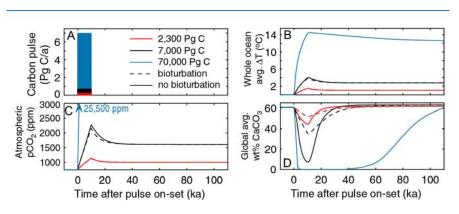


Figure 2-3: Timeseries plots of model output.

Our model does not predict substantial reversal in circulation or weakening of Southern source water production in response to PETM warming, contrary to the interpretation of Nunes and Norris (2006), but in agreement with Bice and Marotzke (2002).

Maps of CaCO₃ wt% results from immediately after the 10 ka pulse (Fig. 2-4) and regional plots of CaCO₃ wt% versus depth (Fig. 2-5) show that no one scenario can

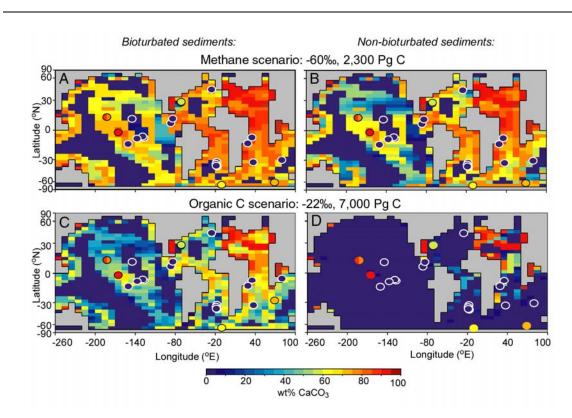


Figure 2-4: Model PETM CaCO₃ wt% with observations superimposed (circular markers). The 70,000 Pg C case is not shown in this figure because there is no CaCO₃ left in the sediments. See Appendix A for data values and sources.

reproduce all of the features of the PETM $CaCO_3$ wt% distribution. In the central equatorial Pacific region, data suggest a CCD at < 3,000 m depth, but sediments from the slightly shallower sites at ~ 2,500 m depth (from Shatsky Rise) retain > 75 wt% $CaCO_3$.

At Walvis Ridge the CCD appears < 1,500 m deep, resulting in 0 wt% CaCO₃ at all Walvis Ridge sites during the PETM. The model, however, displays approximately the same response in both basins. Therefore, the > 75 wt% CaCO₃ in the Pacific can be reproduced by the methane scenarios (Fig. 2-5A), but only at the cost of substantially over-estimating the CaCO₃ content deeper in the Pacific and at Walvis Ridge (Fig. 2-5A,B; we note, however, that our initial scenario underestimates the CaCO₃ wt% at the shallow Pacific sites by 10%). Alternatively, a 7,000 Pg C pulse in the absence of

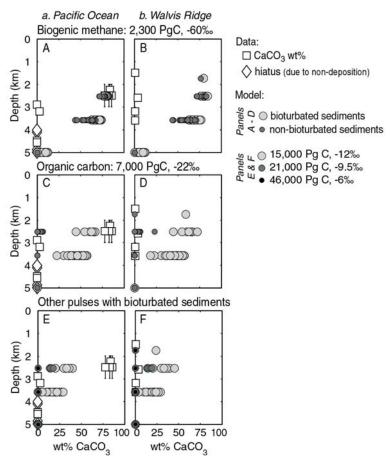


Figure 2-5: PETM model CaCO₃ wt% plotted versus depth with observations superimposed. Results are plotted for data-rich regions of the central equatorial Pacific Ocean (left column) and Walvis Ridge (right column; regions a and b, respectively in Fig. 2-1A).

bioturbation will reproduce the 0 wt% CaCO₃ at depths >3,000 m in the Pacific, and at Walvis Ridge, but only at the cost of losing all CaCO₃ at the 2,500 m sites in the Pacific (Fig. 2-5C,D).

This trade-off is apparent in the distribution of RMSE (Fig. 2-6). For the Pacific region the RMSE is lowest for the 15,000 Pg C pulse with bioturbation (Fig. 2-6A), striking a compromise with 0 wt% CaCO₃ at the 3,600 m model layer, and 20 - 50 wt% CaCO₃ at the 2,500 m model layer. For Walvis Ridge any scenario that dissolves all CaCO₃ will have RMSE of 0. Three scenarios satisfy this requirement: 7,000 Pg C without bioturbation, and 46,000 and 70,000 Pg C with bioturbation (Fig. 2-6B). The sum of the Pacific and Walvis Ridge distributions yields two minima: one for 7,000 Pg C without bioturbation, and the other for 21,000 Pg C with bioturbation (Fig. 2-6C).

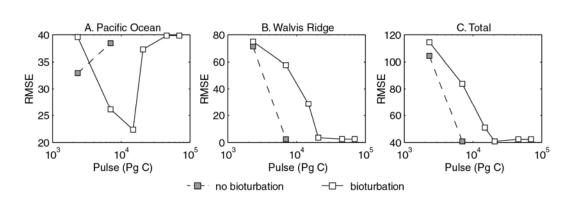


Figure 2-6: Evaluation of model outcomes by comparison of root mean squared errors (RMSE).

One way to reconcile these results is to allow for a decrease in the strength of bioturbation during the PETM at Walvis Ridge and at the deeper Pacific sites. In general, evidence appears to suggest that this was the case at Sites 1220 and 1221 in the Pacific (Shipboard Scientific Party, 2002; Murphy et al., 2006) at paleodepths of 2,900 and

3,200 m, respectively (Rea and Lyle, 2005), and at Walvis Ridge (Shipboard Scientific Party, 2002, 2004b, a). In contrast, sediments at the shallower Pacific Shatsky Rise sites have been reported as "strongly bioturbated" (Colosimo et al., 2006).

What caused the PETM?

Allowing that spatiotemporal variations in the rate of bioturbation can account for the very different responses of the central equatorial Pacific and Walvis Ridge sediments during the PETM, 7,000 Pg C is the smallest pulse modeled here that dissolves enough CaCO₃ to reasonably reproduce observations. The 2,300 Pg C pulse of biogenic methane causes only a small decrease in CaCO₃ wt%, whether bioturbation is present or not. This result is in agreement with Zachos et al. (2005), who concluded that a pulse derived solely from biogenic methane would dissolve insufficient CaCO₃. Our result also offers an alternative to the very high climate sensitivity to CO₂ doubling, 6.8 to 7.8°C, estimated to be necessary to generate PETM warming and a -3 to -5‰ δ¹³C excursion with biogenic methane (Pagani et al., 2006). If 7,000 Pg C is the lower bound on pulse size (given that 7,000 Pg C without bioturbation is the smallest pulse to simulate complete dissolution in the Atlantic, but 15,000 Pg C is the best fit to the Pacific data), then the upper bound on climate sensivity is reduced to 4°C (Fig. 2 of Pagani et al., 2006). Pulses of 21,000 Pg C and greater decrease CaCO₃ wt% too much in our model; the climate sensitivity to CO₂ doubling corresponding to an upper bound on pulse size of 21,000 Pg C is 2.5°C.

A pulse of 7,000 Pg C with δ^{13} C of -22‰ is generally consistent with the hypotheses that organic carbon was released to the atmosphere by peat fires (Kurtz et al., 2003), or by weathering of fossil carbon (Higgins and Schrag, 2006). Alternatively, it could reflect a combination of volcanic CO₂ (i.e., mantle-derived CO₂ with δ^{13} C of -5‰) and thermogenic methane (δ^{13} C of -35 to -50‰, Hunt, 1996) evolved during the emplacement of the North Atlantic Igneous Province. (NAIP; Svensen et al., 2004; Storey et al., 2007). The timing of the PETM appears to coincide with massive eruptions of flood basalts and the emplacement of sill complexes associated with the rifting between Greenland and Europe (Svensen et al., 2004; Storey et al., 2007). NAIP formation is estimated to have yielded between 1,500 and 15,000 Pg C of thermogenic methane (Svensen et al., 2004), so it is not entirely unreasonable that 2,600 to 4,000 Pg C might have been derived from thermogenic methane (with δ^{13} C of -35 and -50‰, respectively), and mixed with higher δ^{13} C CO₂ from basalt degassing.

NAIP volcanism may turn out to be a very economical explanation for the events of the Late Paleocene and the Eocene, accounting for long-term warming, possibly accounting for the PETM, and given that melt production lasted well into the Eocene (Storey et al., 2007), intervals of relatively high rates of melt production may account for PETM-like Eocene hyperthermal events (e.g., Lourens et al., 2005). The NAIP volcanism hypothesis offers the distinct advantage of being relatively easy to test by dating and correlation of volcanic ash layers (e.g., Storey et al., 2007).

Conclusion

In our simulations of the Paleocene-Eocene Earth system, a pulse of 7,000 Pg C with δ^{13} C of -22‰ comes closest to reproducing the changes in sea floor CaCO₃ distribution that took place during the PETM, but it is necessary to permit differences in the rate of bioturbation of Pacific and Atlantic Ocean sediments. A δ^{13} C of -22‰ is consistent with sedimentary organic carbon, or a combination of volcanic degassing and thermogenic methane.

Acknowledgements

This paper represents part of Panchuk's Ph.D. work at Penn State. We wish to thank Robert Marsh and Andrew Yool for setting up the Eocene configuration of the GENIE-1 climate model, and Karen Bice, Jerry Dickens, and an anonymous reviewer for helpful comments on an earlier version of the manuscript. We are very grateful to Tim Bralower, Klaus Keller, and Ray Najjar for their careful reading of the manuscript, detailed written comments, and thoughtful discussions. This work was funded by a Worldwide Universities Network travel grant to Panchuk, Canadian Foundation for Climate and Atmospheric Sciences funding to Ridgwell, and an NSF Biocomplexity grant to Kump.

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Chapter 3

Differences between the carbon isotope history of ocean surface water and the sedimentary record of carbon isotopic excursions during the Paleocene-Eocene Thermal Maximum: A model study²

Abstract

We find that bioturbation and the rate of recovery of the saturation state of seawater are key factors influencing the magitude of the recorded marine $\delta^{13}C$ excursion as well as the record of changes in CaCO₃ wt%. We conclude that, while the marine $\delta^{13}C$ excursion and CaCO₃ wt% records may not be consistently reliable indicators of the amount and $\delta^{13}C$ of the carbon added to the Earth system to cause the PETM, the timescales of the system recovery as indicated by those records might be. Differences in these timescales from one site to another may also provide valuable information about the response of the Earth system to PETM climate change.

Introduction

Mass balance calculations have been used to determine the carbon source that caused the Paleocene-Eocene Thermal Maximum (PETM) event: given the carbon isotopic composition (δ^{13} C) of possible sources, how much carbon from a particular

² K. Panchuk, A. Ridgwell, L. Kump, and T. Bralower, for submission to Paleoceanography.

source would be required to produce the carbon isotopic excursion (CIE) observed in the Latest Paleocene ocean-atmosphere system, and is that quantity consistent with our understanding of the state of the Earth system at that time? A unique solution has not come of this approach because there are many combinations of δ^{13} C and pulse size that could provide the right excursion. Additional constraints may be brought to bear on the problem, such as spatial and temporal changes in the distribution of CaCO₃ sediments (Chapter 2). However, uncertainty regarding the magnitude of the PETM CIE itself remains to be addressed (e.g., Pagani et al., 2006).

There is a puzzling discrepancy in the magnitude of the PETM CIE recorded in terrestrial versus marine environments. On land the CIE is -5 to -7‰ (Bowen et al., 2004 and references therein), whereas foraminifera from marine sediments display average excursions of -3 to -4‰ (Table 3-1). The terrestrial δ¹³C excursions may have been amplified relative to the marine by increased turnover of organic matter in soils and increased isotopic discrimination during photosynthesis under conditions of higher relative humidity (Bowen et al., 2004). However, this does not explain the range of excursions in the marine record, from -1‰ to larger than -4‰ (Table 3-1). CaCO₃ dissolution might account for this variability; dissolution and/or lack of deposition of sediments is characteristic of the PETM, although CaCO₃ wt% varies widely, with relatively little change in the CaCO₃ content of sediments at some locations (e.g., Shatsky Rise in the Pacific Ocean, Colosimo et al., 2006), but complete dissolution of CaCO₃ at others (e.g., Walvis Ridge in the Atlantic, Zachos et al., 2005).

Table 3-1: δ^{13} C values of planktonic foraminifera and bulk CaCO₃

Location	Material	δ^{13} C		Course
		baseline	excursion	- Source
Pacific Ocean				
Kerguelen Plateau (Site 738)	A. soldadoensis	3.8	-4.0	Lu and Keller (1993)
Shatsky Rise				
1209	M. velascoensis	4.0	-3.2	Zachos et al. (2003)
	A. soldadoensis	4.0	-3.7	Zachos et al. (2003)
	bulk CaCO ₃	2.8	-2.5	Colosimo et al. (2006)
1210	bulk CaCO ₃	2.6	-2.1	Colosimo et al. (2006)
1211	bulk CaCO ₃	2.9	-2.7	Colosimo et al. (2006)
1212	bulk CaCO ₃	3.2	-2.7	Colosimo et al. (2006)
Equatorial Pacific				
1220	bulk CaCO ₃	1.0	-0.6	Nunes and Norris (2005)
1221	bulk CaCO ₃	1.9	-0.9	Nunes and Norris (2005)
Caribbean Sea				
Nicaraguan Rise (Site 1001)	bulk CaCO ₃	2.0	-3.0	Bralower et al. (1997)
Atlantic Ocean				
Bay of Biscay (Site 401)	M. subbotinae	4.3	-2.0	Pardo et al. (1997)
Goban Spur (Site 549)	M. subbotinae	4.7	-4.5	Stott et al. (1996)
Blake Nose (Site 1051)	bulk CaCO ₃	2.2	-1.6	Bains et al. (1999)
Walvis Ridge				
525	bulk CaCO ₃	2.4	-2.5	Thomas and Shackleton (1996)
1262	bulk CaCO ₃	2.3	-2.2	Zachos et al. (2005)
1263	bulk CaCO ₃	2.0	-2.9	Zachos et al. (2005)
1265	bulk CaCO ₃	2.2	-2.7	Zachos et al. (2005)
1266	bulk CaCO ₃	2.1	-2.0	Zachos et al. (2005)
1267	bulk CaCO ₃	2.2	-2.6	Zachos et al. (2005)
Southern Ocean				
Maud Rise				
689	A. mckannai	2.9	-2.9	Thomas and Shackleton (1996)
690	Acarinina	3.2	-4.0	Thomas et al. (2002)
	bulk CaCO ₃	2.9	-2.8	Bains et al. (1999)

If smaller δ^{13} C excursions reflect gaps in the sedimentary record, we would expect to see that smaller excursions occur at sites with greater decreases in CaCO₃, but this is not obvious in the published data (Fig. 3-1; locations are shown in Fig. 3-2). In Fig. 3-1A, the CIE magnitudes in planktonic foraminiferal records and bulk carbonate are plotted against the decrease in CaCO₃ wt% normalized to pre-PETM values, Δ_n CaCO₃:

$$\Delta_n CaCO_3 = \frac{CaCO_{3,initial} - CaCO_{3,final}}{CaCO_{3,initial}}$$
(Eq. 3-1)

where $CaCO_{3,initial}$ and $CaCO_{3,final}$ are the $CaCO_3$ wt% for a site before the PETM and at the peak of the PETM, respectively. For sediments with $CaCO_3$ prior to the PETM, but none during the PETM, $\Delta_n CaCO_3 = 1$. The advantage of this formulation is that it distinguishes between cases with the same change in $CaCO_3$ wt%, but different initial $CaCO_3$ wt%. In other words, sites where $CaCO_3$ wt% decreases from 80% to 60% and from 50% to 30% both have decreases of 20 wt% $CaCO_3$, but more $CaCO_3$ is dissolved in the former case than in the latter.

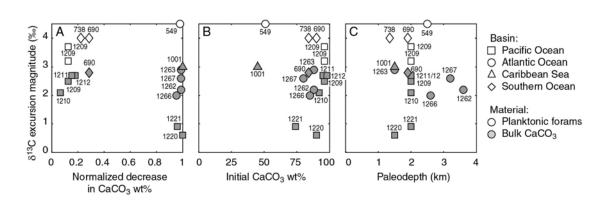


Figure 3-1: δ^{13} C excursion magnitude in planktonic foraminifera and bulk CaCO₃ versus: A) the decrease in CaCO₃ wt% normalized to initial values; B) initial CaCO₃ wt%; C) paleodepth. Data are labeled with DSDP/ ODP site numbers. See Appendix A for data and sources.

In general, more sites with larger excursions do show a smaller decrease in CaCO₃, but there is remarkably little difference between the bulk CaCO₃ CIE magnitudes at Walvis Ridge (Sites 1262-1267), where $\Delta_n CaCO_3 \sim 1$, and those at Shatsky Rise (Sites 1209-1212) and Maud Rise (Site 690) where $\Delta_n CaCO_3 < 0.3$. Another interesting feature of Fig. 3-1A is that, for $\Delta_n CaCO_3$ up to ~ 0.3 , CIE magnitudes appear to increase rather than decrease as $\Delta_n CaCO_3$ increases. At Site 1209 the excursion magnitudes in planktonic foraminifera are 3.2 and 3.7% with $\Delta_n CaCO_3 \sim 0.1$, whereas Sites 738 and 690 have 4% excursions and $\Delta_n CaCO_3$ of 0.2 to 0.3. Similarly, at Site 1210 the bulk $CaCO_3$ CIE magnitude is 2.1% and $\Delta_n CaCO_3$ is ~ 0.1 , but at Site 690 the bulk $CaCO_3$ CIE magnitude is 2.8% with $\Delta_n CaCO_3$ of 0.3. The excursion magnitudes do not appear to be a function of initial $CaCO_3$ wt% (Fig. 3-1B) or paleodepth (Fig. 3-1C).

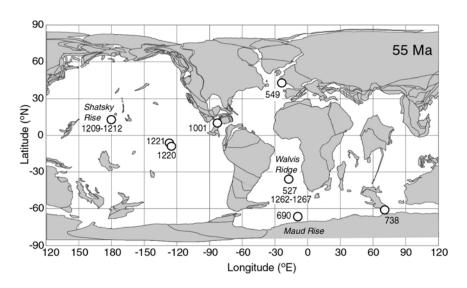


Figure 3-2: Late Paleocene paleogeography showing the location of sites in Fig. 3-1. Paleogeographic reconstruction by the Advanced Plate Reconstruction Service of the Ocean Drilling Stratigraphic Network, www.odsn.de.

The relationship between the magnitude of the CIE and CaCO₃ content in Fig. 3-1A does not rule out the possibility that CaCO₃ dissolution muted CIE records in deep sea sediments, but it does suggest that other factors may have played a role. Our objective in this paper is to determine what those factors might have been and to address the underlying question, which is broadly applicable to events throughout Earth history, "How much different is the δ^{13} C record preserved in sediments from the history of δ^{13} C in the seawater from which those sediments originated?"

Of course, this question is difficult to answer because we have no direct samples of ancient seawater with which to reconstruct the oceanic δ^{13} C history. Instead, we have approached the problem with a numerical Earth system model, GENIE-1, which allows us to analyze the propagation of a simulated PETM δ^{13} C signal in both seawater and preserved sediments. GENIE-1 has a 3D-ocean biogeochemistry model, as well as a sediment component that both simulates CaCO₃ preservation on the sea floor and generates a stratigraphic record as sediments accumulate. These model features allow us to compare the PETM δ^{13} C signal as it appears in surface ocean dissolved inorganic carbon (DIC) to its expression in newly produced CaCO₃, in core-top CaCO₃, and in the stratigraphic record.

Our previous use of GENIE-1 to study the PETM (Chapter 2) focused on calculating the spatial distribution of CaCO₃ wt% to provide a constraint on the source of C necessary to generate the presumed PETM δ^{13} C excursion. Here we extend that work to explore the relationship between actual and preserved carbon isotope excursions to both further constrain the sources and sizes of C inputs during the PETM, and more

generally assess the nature of the isotopic record of C-cycle variations in the geologic past.

We explore several factors that may have influenced the $\delta^{13}C$ excursion. These include the potential for fractionation-related damping of the excursion in newly produced CaCO₃, the effects of spatial variability in the pre-PETM $\delta^{13}C$ of surface ocean DIC ($\delta^{13}C_{DIC}$), the effects of dissolution of "old" CaCO₃ with relatively high $\delta^{13}C$ on the CIE, and the role of bioturbation. We find that the primary determinant of the CIE magnitude preserved in sediments is the delay between the onset of isotopic recovery of surface water DIC and the resumption of sedimentation at a particular model site.

Methods

The GENIE-1 Earth system model (Edwards and Marsh, 2005) is driven by a frictional geostrophic 3-D ocean circulation model (Edwards, 1996) coupled to a 2-D energy-moisture balance model (EMBM) of the atmosphere and a dynamic-thermodynamic sea-ice model based on that of Weaver et al. (2001). Also included is a marine carbon-cycle model (Ridgwell et al., 2007) and a model of biogenic carbonate sedimentation (Ridgwell and Hargreaves, 2007). Configuration of the model for the PETM is discussed in Chapter 2. Detailed descriptions of all components are available in the sources cited above, so here we focus on the model components of most interest to this investigation, those that deal with the production and preservation of CaCO₃ and isotopic changes in the carbon cycle. These components are illustrated schematically in Fig. 3-3.

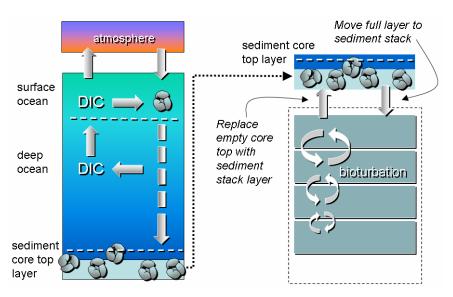


Figure 3-3: Schematic illustration of the GENIE-1 Earth system model. White arrows indicate transport directions. Forams (\mathfrak{S}) symbolize newly produced CaCO₃ with δ^{13} C representative of the surface ocean.

In the model, production of new CaCO₃ in the surface ocean is ultimately determined by the PO_4^{3-} concentration in the surface layer. Particulate organic carbon (POC) export is calculated from the PO_4^{3-} concentration using a Monod growth formulation. Export of CaCO₃ is determined by a prescribed rain ratio (see Chapter 2), which sets the CaCO₃ export relative to the POC export. New CaCO₃ has δ^{13} C equal to that of surface ocean DIC plus a positive fractionation term that increases with temperature (after Mook, 1986). New CaCO₃ and POC fall through the water column and undergo dissolution according to an exponential function with length scales of 1055 m and 556 m, respectively (Ridgwell et al., 2007).

All POC reaching the sea floor is remineralized and returned to the water column, and the CO₂ evolved is available to dissolve CaCO₃ within the sediments. CaCO₃ reaching the sea floor is added to the core-top sediment layer with 1 cm maximum

thickness. If the core-top layer is full, it is moved down into a sediment stack, and accumulation on the sea floor continues until sediments in the core-top layer are once again 1 cm thick. If the sea floor sediment layer is emptied, as by dissolution, then the top sediment stack layer is moved up to take its place. In this way the model sediment core may have gaps if long periods of non-deposition occur, or if layers of the core are "unburied" and dissolved, similar to the "burn-down" thought to take place at the PETM (e.g., Zachos et al., 2005; Colosimo et al., 2006). Non-carbonate detrital material is accumulated along with CaCO₃ at 0.18 g cm⁻² ka⁻¹. This rate was tuned to improve the model approximation of the pre-PETM CaCO₃ distribution, but it differs little from a value typical of the modern Pacific Ocean and one used in previous long-term fossil-fuel CO₂ removal simulations by Archer et al. (1998), 0.16 g cm⁻² ka⁻¹. Without the non-carbonate detrital component, the CaCO₃ wt% could be only 0% or 100%.

Within the sediment stack, sediments can undergo bioturbation. Bioturbation is represented as mixing between adjacent sediment layers, where the mixing rate is a maximum of 16 cm² ka⁻¹ in the top layer and decreases with an e-folding depth of 1 cm (Ridgwell, 2001). Bioturbation is important because it can bring CaCO₃ into the core-top layer as other CaCO₃ is dissolved and therefore increase the amount of CaCO₃ available to react with CO₂ in seawater. This ultimately increases the CO₂ uptake capacity of the ocean.

Although model time may be applied directly to DIC or core-top records, this does not work for the sediment stack because any one layer may contain sediments accumulated at different model times. This is accommodated by introducing an age tracer. The age tracer behaves as any other conservative tracer, and the age of a sediment

stack layer is the weighted average of the ages of sediments making it up, in the same way that the $\delta^{13}C$ of bulk carbonate represents the weighted average of $\delta^{13}C$ of its components.

The simulations shown here include three of the scenarios that have been suggested to explain the PETM: release of biogenic methane from methane clathrates (δ^{13} C of -60%), oxidation of organic carbon (δ^{13} C of -22%), and emission of mantle CO₂ by volcanism (δ^{13} C of -5%). The size of the CO₂ pulse in each case is that which generates a -4% excursion in the combined ocean-atmosphere carbon reservoir. For the biogenic methane, organic carbon, and mantle CO₂ scenarios, the necessary pulse sizes are 2,300, 7,000, and 70,000 Pg C, respectively. Three additional simulations were run to fill the gap between the 7,000 and 70,000 Pg C pulses. These have δ^{13} C of -12, -9.5, and -6% and pulse sizes of 15,000, 21,000, and 46,000 Pg C, respectively. These scenarios can be viewed as some of the possible combinations of thermogenic methane (δ^{13} C of -35 to -50%, Hunt, 1996) and volcanic (i.e., mantle-derived) CO₂ discharged into the atmosphere during North Atlantic Igneous Province (NAIP) volcanism (e.g., Svensen et al., 2004; Storey et al., 2007). Indeed, pulses with δ^{13} C of -22% and lower could also be consistent with carbon derived from NAIP volcanism.

The model was spun up for 150 ka with pCO₂ of 750 ppm and a flux of 50 Tmol HCO_3^- a⁻¹ from rivers (see Chapter 2 for details) to achieve geochemical steady state and to build up a sufficient depth of sediments. After spin-up a pulse of carbon in the form of CO_2 was released to the atmosphere at a uniform rate over 10,000 years (after Dickens, 2003). At the end of 10,000 years the pulse was stopped and the model was run

for an additional 100,000 years. Each scenario was run with bioturbation, and the 2,300 and 7,000 Pg C scenarios were also run without bioturbation. Here we examine the sedimentary records produced by the model to understand how the record of $\delta^{13}C_{DIC}$ preserved in sediments might have been affected by dissolution.

Results

The response of the model climate to the ensemble of carbon-release scenarios is summarized in Fig. 3-4B,C: Atmospheric pCO₂ ranged between 1,140 (2,300 Pg C pulse) and 25,500 ppm (70,000 Pg C pulse) immediately after 10 ka of carbon addition (Fig 3-4B). At this time sea surface temperature (SST) was also at its peak, having warmed by 1.5°C to 14°C (Fig. 3-4C). A doubling of atmospheric pCO₂ (from 750 ppm to 1500 ppm) yielded 1°C of warming of the globally averaged SST (Fig. 3-4D).

There are several key features in the timeseries of $\delta^{13}C$ _{DIC}, core-top layer CaCO₃ $\delta^{13}C$ ($\delta^{13}C_{CT}$), and core-top layer CaCO₃ wt% (Figs. 3-4E-G, respectively) that provide important clues about the factors affecting how the pulse of low $\delta^{13}C$ CO₂ to the atmosphere is eventually recorded in the sediments.

Surface ocean DIC δ^{13} C

The larger the CO_2 pulse to the atmosphere, the smaller the $\delta^{13}C$ $_{DIC}$ excursion (Fig. 3-4E). The CIE is the largest for the 2,300 Pg C case, at nearly -4‰, and the smallest for the 70,000 Pg C case, at ~-2‰. The decrease in CIE magnitude reflects a

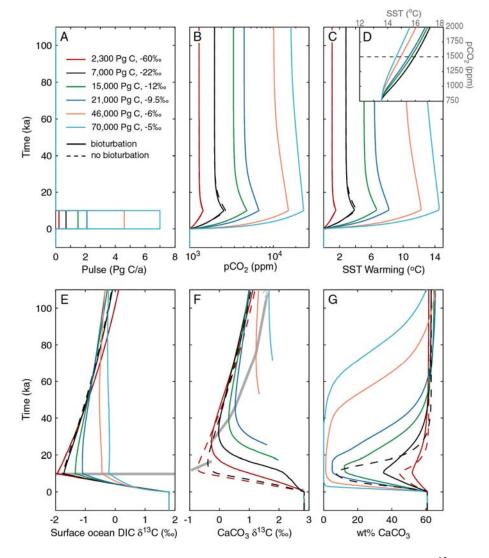


Figure 3-4: Globally averaged timeseries of system properties. A) low $\delta^{13}C$ pulse input signal; B) atmospheric pCO₂; C) increase in sea surface temperatures (SST); D) sensitivity of global average SST to increasing atmospheric CO₂ (dashed line marks pCO₂ doubling from the initial 750 ppm); E) $\delta^{13}C$ of surface ocean DIC (grey line indicates the peak of the excursion at 10 ka); F) $\delta^{13}C$ of CaCO₃ in sediment core-top layer (grey line intersects results at peak excursion); G) CaCO₃ content of sediment core-top layer.

decrease in the net burial of CaCO₃ as seawater saturation with respect to calcite (Ω) decreases, and eventually seawater becomes undersaturated (Ω <1; Ω = 1 indicates saturation, and Ω >1 indicates supersaturation). The flux of carbon from rivers is no

longer balanced by the flux of carbon to the sediments. Furthermore, old CaCO₃, deposited prior to the onset of the PETM, is dissolved. Because the old CaCO₃ was produced from surface ocean DIC when $\delta^{13}C_{DIC}$ was higher, dissolution of old CaCO₃ contributes higher $\delta^{13}C$ carbon to the DIC pool, offsetting some of the isotopic effects of the low $\delta^{13}C$ carbon pulse. For the 2,300 Pg C cases, the system carbon reservoir is $\sim 3,300$ Pg C larger by the end of the 10 ka carbon pulse, indicating that $\sim 1,000$ Pg C was added to the system above the 2,300 Pg C pulse. For the 7,000 Pg C cases the system carbon reservoir is > 2,000 Pg C larger than expected, and for the 70,000 Pg C case decreased net deposition accounts for a further 7,600 Pg C.

There is an additional level of complexity to the $\delta^{13}C_{DIC}$ response: for a given carbon-release scenario, DIC exhibits a range of $\delta^{13}C$ values, and the range is wider for larger carbon pulses (Fig. 3-5A). For the 2,300 Pg C cases, the range is $\sim 0.5\%$, and for the 70,000 Pg C case the range is nearly 2‰. The range varies smoothly with the initial $\delta^{13}C_{DIC}$ (Fig. 3-5B).

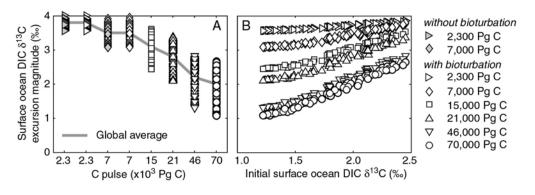


Figure 3-5: A) $\delta^{13}C_{DIC}$ CIE magnitudes of individual model sites for each carbon-release scenario. B) Relationship between the initial $\delta^{13}C_{DIC}$ and the $\delta^{13}C_{DIC}$ excursion magnitude, reflecting two-component mixing.

This relationship is an expected outcome of the two-component mixing taking place between the newly-added low- δ^{13} C CO₂ and surface ocean DIC. By re-arranging mass balance and isotope mass balance equations (Eqs. 3-2 and 3-3, respectively), where M is mass, δ is δ^{13} C, and the subscripts i, a, and f refer to initial, added, and final quantities, respectively, we see that the δ^{13} C_{DIC} excursion ($\varepsilon = \delta_f - \delta_i$) is proportional to the difference between the δ^{13} C of the CO₂ pulse and the initial δ^{13} C_{DIC}, and to the ratio of the mass of added carbon to the final (total) mass of carbon in the system:

$$M_f = M_i + M_a$$
 (Eq. 3-2)

$$\delta_f M_f = \delta_i M_i + \delta_a M_a \tag{Eq. 3-3}$$

$$\varepsilon = \left(\delta_a - \delta_i\right) \frac{M_a}{M_f} \tag{Eq.3-4}$$

The greater the difference in δ^{13} C between the initial reservoir and the added carbon, the larger the excursion magnitude will be. The larger the ratio of M_a to M_f , the greater the range of excursions. Consequently, the magnitude of the excursion due to biogenic methane is large because of the large difference between the δ^{13} C of methane (-60‰) and the δ^{13} C of surface ocean DIC (~-1 to -2‰), but the range of excursion magnitudes is small. On the other hand, the excursion caused by mantle CO₂ (-5‰) is much smaller, but the larger pulse size amplifies the effects of variations in the initial δ^{13} C_{DIC}.

Core-top CaCO₃ record

Like the globally averaged $\delta^{13}C_{DIC}$, the magnitude of the CIE in $\delta^{13}C_{CT}$ also decreases for larger pulse sizes, but the magnitudes of CaCO₃ CIEs are smaller than those

of $\delta^{13}C_{DIC}$ (Fig. 3-4F). In general, the global average $\delta^{13}C_{DIC}$ is the upper bound of the range of CIE magnitudes in the core-top record (Fig. 3-6A).

Some of the difference between DIC and core top $CaCO_3$ excursion magnitudes can be accounted for by the increase in carbon isotopic fractionation between surface ocean DIC and newly produced $CaCO_3$. Initially, fractionation results in $CaCO_3$ $\delta^{13}C$ values of ~1‰ higher than the DIC from which the $CaCO_3$ is formed. As the temperature increases in response to the CO_2 pulse, the positive fractionation becomes larger, and the net increase in fractionation offsets the negative $\delta^{13}C$ excursion. The amount of damping by changes in fractionation (Fig. 3-6C) corresponds roughly to the difference between the $\delta^{13}C_{CT}$ and $\delta^{13}C_{DIC}$ excursions in the case of the three or four largest pulses, but for the smaller pulses it leaves most of the difference unaccounted for.

Another feature of the globally-averaged $\delta^{13}C_{CT}$ timeseries is that, whereas the peak of the CIE in the surface ocean occurs at the end of the 10 ka carbon pulse for all cases, the peak of the CIE in the core-top $CaCO_3$ is delayed, and the delay is longer for larger carbon pulses (heavy grey lines intersect the peak $\delta^{13}C$ values in Fig. 3-4F, cf. grey line in Fig. 3-4E). The time elapsed between the $\delta^{13}C_{DIC}$ peak and the $\delta^{13}C_{CT}$ peak (Δt_p and Δt_p ' in Fig. 3-7A, where the prime indicates the presence of bioturbation) is due in part to a delay between the onset of the $\delta^{13}C_{DIC}$ (i.e., at the start of the carbon pulse) and the onset of the $\delta^{13}C_{CT}$ CIE (the intervals Δt_o and Δt_o ' in Fig. 3-7A). The delay in the onset of the $\delta^{13}C_{CT}$ CIE reflects the interval after the pulse begins during which seawater becomes undersaturated with respect to calcite, so there is no net deposition of $CaCO_3$ (Δt_o and Δt_o ' in Fig. 3-7B) and therefore no medium to record $\delta^{13}C_{DIC}$.

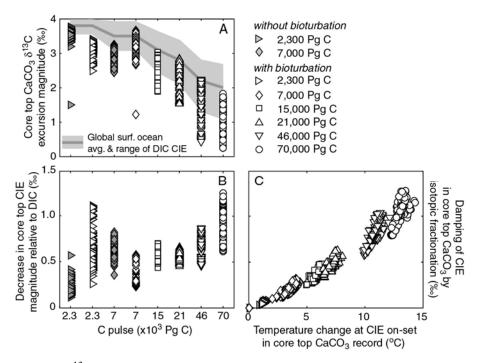


Figure 3-6: A) $\delta^{13}C_{CT}$ excursion magnitude for each carbon-release scenario. B) Damping of the $\delta^{13}C_{CT}$ CIE relative to the $\delta^{13}C_{DIC}$ CIE for each scenario. C) Damping of the $\delta^{13}C_{CT}$ CIE due to temperature-related changes in isotopic fractionation between DIC and CaCO₃.

For the case without bioturbation illustrated in Fig. 3-7, a small part of the beginning of the excursion is recorded while net sedimentation persists (i.e., while Ω >1, up to a in Fig. 3-7A,B). When all available CaCO₃ is used up, the rate of dissolution is exactly equal to the rate of deposition, and the record of the excursion stops (c in Fig. 3-7B,C). Approximately 4 ka after the carbon pulse stops, seawater at this site is no longer undersaturated with respect to calcite. The rate of dissolution decreases below the rate of deposition, CaCO₃ begins to accumulate, and the record of δ ¹³C_{DIC} resumes, having missed the peak of the excursion.

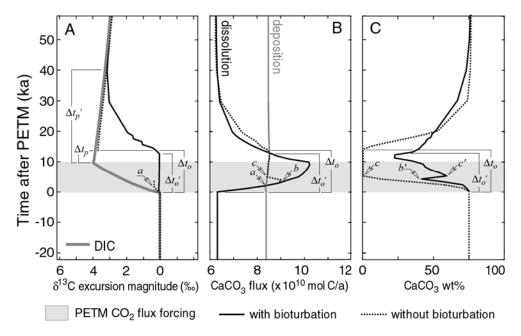


Figure 3-7: Model DIC and core-top timeseries for the 7,000 Pg C cases from Shatsky Rise. All points labelled with the same letter occur at the same time. Labels with primes indicate the case with bioturbation. The intervals Δt_p and $\Delta t_p'$ are the delays between the peak of the excursion in DIC and the peak of the excursion in the core-top record. The intervals Δt_o and $\Delta t_o'$ are the delays between the onset of the excursion in DIC and the onset of the excursion in the core-top record. A) $\delta^{13}C_{CT}$ excursion magnitude superimposed on the $\delta^{13}C_{DIC}$ record (heavy grey line). B) Rates of dissolution (black lines) and deposition (grey line; deposition is the same for bioturbated and non-bioturbated cases). C) Core-top CaCO₃ wt% timeseries.

When bioturbation is present, however, the timing of the peak is affected by the on-going supply of old $CaCO_3$ from the sediments in addition to the length of time required for recovery of Ω . At the onset of the pulse, the dissolution flux follows the same trajectory as before, until the point where the non-bioturbation scenario ran out of available $CaCO_3$ (Fig. 3-7B, b). Here the dissolution flux for the bioturbation case continues to increase until the pulse stops, because Ω is still decreasing as a result of CO_2 addition, and because bioturbation provides a ready supply of old $CaCO_3$ to be dissolved.

In this case, CaCO₃ wt% does not go to zero because dissolution is sufficiently slow compared to the resupply of old CaCO₃ by bioturbation. In fact, the difference in rates between these two processes results in double CaCO₃ wt% minima. The first minimum (Fig. 3-7C, *b'*) occurs when the old CaCO₃ immediately available in the surface layer is used up. Despite the fact that the rate of dissolution continues to rise in this case, there is a brief period during which the CaCO₃ wt% increases because the rate of bioturbation is supplying CaCO₃ in excess of the dissolution rate. When the dissolution rate increases beyond the sum of the supply rates of bioturbation and deposition (Fig. 3-7C, *c'*), CaCO₃ wt% falls again until shortly after the pulse stops.

When the rate of dissolution falls below the rate of accumulation, the $\delta^{13}C$ excursion appears in the core-top record. Until then, $\delta^{13}C_{CT}$ continues to reflect the $\delta^{13}C$ of old CaCO₃. The time elapsed between the onset of the $\delta^{13}C$ excursion in bioturbated sediments and its peak is a function of the rate of new CaCO₃ accumulation relative to the rate of mixing of new CaCO₃ with old. The excursion peaks and matches the $\delta^{13}C_{DIC}$ trajectory when enough new CaCO₃ has accumulated such that dilution by old, higher $\delta^{13}C$ CaCO₃ becomes negligible. The maximum possible $\delta^{13}C_{CT}$ excursion in either case is therefore dependent upon the rates of recovery of sedimentation and of $\delta^{13}C_{DIC}$.

It appears that the presence or absence of bioturbation has substantial consequences for the preservation of $CaCO_3$ in the sediments and on the $\delta^{13}C_{CT}$ profile. Where sediments are not bioturbated, $\delta^{13}C$ and $CaCO_3$ wt% decrease more rapidly. The $\delta^{13}C_{CT}$ CIE is larger and sharper in cases without bioturbation, and the global average $CaCO_3$ wt% is lower compared to sediments with bioturbation for the same pulse size

(Figs. 3-4G, 3-7C). Given the response of core-top CaCO₃ content, it is perhaps counterintuitive that *less* CaCO₃ is dissolved in the absence of bioturbation than when bioturbation is present. Without bioturbation, the CaCO₃ available for dissolution is limited by what is immediately available at the sediment surface. When this is depleted, a cap of non-carbonate detrital material remains and impedes interaction by diffusion with CaCO₃ deeper in the sediment column. In contrast, with bioturbation previously-buried CaCO₃ is moved back into the sediment surface layer and becomes available for dissolution. Consequently, the CO₂ uptake capacity of the ocean is reduced in the absence of bioturbation. For the 7,000 Pg C case without bioturbation, both pCO₂ and temperature are slightly higher than the case with bioturbation (Fig. 3-4B,C). This effect would be more pronounced for larger carbon pulses because the 7,000 Pg C case without bioturbation already comes close to dissolving most of the available CaCO₃ globally (see Fig. 2-4D), so more of the carbon from larger pulses would remain in the atmosphere.

Sediment-core CaCO₃ record

The difference between the core-top record and the sediment-core record is that the core top is an instantaneous snapshot of the sediment conditions on the sea floor, whereas the sediment core reflects the processes that occur after the surface layer becomes buried. This difference does not appear to affect the magnitude of the δ^{13} C excursion that is recorded; for the most part there is a one-to-one relationship between core top and sediment-core δ^{13} C excursion magnitudes (Fig. 3-8A).

The delay between the peak of the $\delta^{13}C_{DIC}$ excursion and the $\delta^{13}C$ excursion in CaCO₃ tends to be smaller in the sediment core than in the core top, except for the 70,000 Pg C case, and some model sites in the 46,000 Pg C case (Fig. 3-8B). Variability in delay times is also slightly less for the sediment core, although the delay times in the sediment cores tend to fall within the range of delay times for the core-top records. The delay in the core-top record depends in part on when net deposition resumes. This is also true for the sediment-core record, but the timing of the excursion can then be modified by bioturbation within the sediment stack; while bioturbation delays the peak of the excursion in the core-top record, it advances the peak in the sediment core relative to the core-top excursion, and advances the onset of the excursion so that it appears to begin

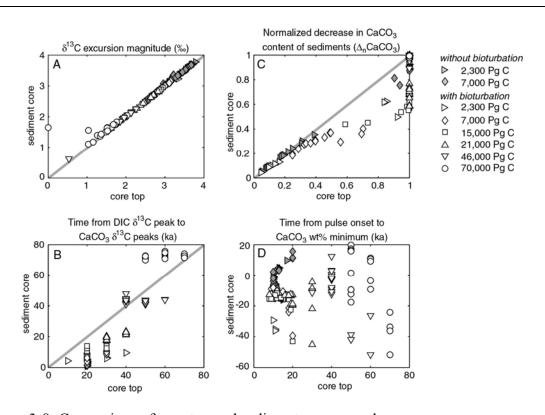


Figure 3-8: Comparison of core top and sediment-core records.

before the $\delta^{13}C_{DIC}$ excursion (Fig. 3-9A). This occurs because dissolution of CaCO₃ in sediments eventually results in new CaCO₃ being deposited on sediments older than the PETM, and bioturbation mixes the $\delta^{13}C$ of the new CaCO₃ downcore. In contrast, when bioturbation is not present, the magnitude and timing of the excursion in the core top and the sediment core are essentially the same (Fig. 3-9B).

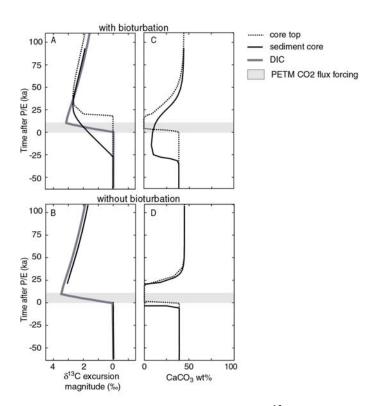


Figure 3-9: Comparison of core top and sediment-core $\delta^{13}C$ excursion and CaCO₃ wt% profiles from a 5000 m deep site at Walvis Ridge for the 7,000 Pg C cases with and without bioturbation.

The change in CaCO₃ wt% during the PETM is very similar for sediment core and core-top records when the decrease in CaCO₃ wt% is relatively small, but when larger fractions of CaCO₃ in sediment dissolves, there is less of a decrease in the sediment-core CaCO₃ content than in the core-top CaCO₃ content (Fig. 3-8C). This is also an effect of

bioturbation; even when the $CaCO_3$ wt% of the core-top record is zero, the sediment core may record a non-zero $CaCO_3$ wt% because bioturbation mixes carbonate-free sediments with older sediments still containing $CaCO_3$ (Fig. 3-9C). As with the $\delta^{13}C$ record, downcore mixing of carbonate-free sediment with older $CaCO_3$ -rich sediment advances the onset of the $CaCO_3$ decline so that it appears to occur before the carbon pulse is added. Without bioturbation, 0 wt% $CaCO_3$ in the core top will be recorded as such in the sediment core (Fig 3-9D).

Comparison of observed and modeled CIE magnitudes

The CIE magnitudes observed in the data do not have a straightforward relationship with the decrease in CaCO₃ wt% (Fig. 3-1), and neither do the excursion magnitudes in the sediment core (Fig. 3-10). The 2,300 and 7,000 Pg C cases without bioturbation have the largest excursion magnitudes, ~3.5‰ on average; the 2,300 Pg C case spans a range of Δ_n CaCO₃ from 0.1 to 1 (from a 10% decrease in CaCO₃ wt% to a 100% decrease) with no apparent change in CIE magnitude, and the 7,000 Pg C case plots in the range of Δ_n CaCO₃ from 0.7 to 1. The 2,300 and 7,000 Pg C cases with bioturbation have average CIE magnitudes of ~3 ‰, and also cover a wide range of Δ_n CaCO₃ without showing an obvious decrease in CIE magnitude when the decrease in CaCO₃ wt% is greater. The other cases tend toward higher values of Δ_n CaCO₃ and lower CIE magnitudes as pulse sizes become larger, but they still display a considerable range of excursion magnitudes. For example, the 70,000 Pg C case has excursions between 0 and -2‰ despite having Δ_n CaCO₃ of 1 or close to 1 in all cases.

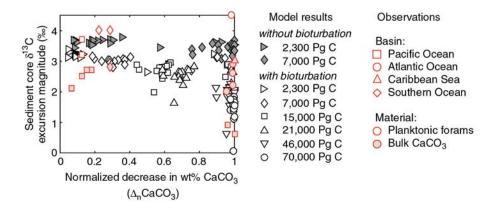


Figure 3-10: Comparison of $\delta^{13}C$ excursion magnitude vs. normalized decrease in wt% CaCO₃ between model sediment-core results and observations. See Appendix A and Table 3-1 for data and sources.

No single case modeled here is consistent with all of the observations. In the low $\Delta_n \text{CaCO}_3$ region of Fig. 3-10 (left-hand side) the Pacific ocean data from planktonic foraminifera fall within the range of results for the 2,300 Pg C cases with and without bioturbation, but the Southern Ocean planktonic foram observations have $\delta^{13}\text{C}$ excursions larger than any of the model results. Excursions recorded in bulk carbonate from the Pacific fall below the range of model results, although the bulk carbonate measurement from the Southern Ocean falls within the range of the 7,000 Pg C case with bioturbation. On the right hand side of the plot, all observations have $\Delta_n \text{CaCO}_3 > 0.9$. The largest excursion is -4.5%, from northern Atlantic Ocean planktonic forams, and well above the range of model results. The Caribbean Sea observation from bulk carbonate falls within the range of the 7,000 to 15,000 Pg C cases; the Atlantic Ocean (Walvis Ridge) bulk carbonate observations fall within the range of the 7,000 to 46,000 Pg C cases; and the two bulk carbonate measurements from the Pacific are within the range of the 46,000 to 70,000 Pg C cases.

Discussion

In this study pulses of CO₂ were added to the atmosphere in amounts calculated to decrease the δ^{13} C of the system carbon reservoir (ocean plus atmosphere) by 4‰. However, at each step toward recording the change in δ^{13} C in the sediments, there are factors at play that can modify the δ^{13} C signal: The δ^{13} C_{DIC} CIE magnitude depends on the initial $\delta^{13}C_{DIC}$, where a greater difference between the initial $\delta^{13}C_{DIC}$ and the $\delta^{13}C$ of the CO₂ pulse yields a larger excursion (Fig. 3-5, Eq. 3-4), and on the amount of old, relatively high δ^{13} C CaCO₃ that is dissolved as a result of the carbon pulse and added back into the ocean-atmosphere reservoir (Fig. 3-4E). The formation of CaCO₃ from DIC includes a temperature-dependent fractionation step that damps the CIE in new CaCO₃ compared to DIC (Fig. 3-6). The CIE magnitude is further decreased if there is a long period of undersaturation of seawater with respect to calcite. While Ω <1, no new $CaCO_3$ accumulates, so there is no record of the $\delta^{13}C$ history in the surface ocean. When the perturbation to Ω is small, deposition recovers more rapidly and more of the $\delta^{13}C_{DIC}$ excursion is recorded. When the perturbation is large and deposition recovers later, the sediment record may miss much of the largest part of the excursion. In the presence of bioturbation, new CaCO₃ is diluted by old, and it is not until enough new CaCO₃ has accumulated to make dilution negligible that sediments will begin to reflect $\delta^{13}C_{DIC}$. Bioturbation can affect sediments long after they are buried. While this appears to have a negligible effect on the magnitude of the δ^{13} C excursion recorded in the sediment core, bioturbation mixes low CaCO₃ or CaCO₃-free sediments with sediments having higher CaCO₃ wt%. Consequently, a higher CaCO₃ wt% is recorded than may have existed.

Overall, this behavior does not result in a straightforward relationship between the magnitude of the δ^{13} C excursion and the gap in the sedimentary record as indicated by the normalized decrease in CaCO₃ wt% (Fig. 3-10), but it does suggest a framework for interpreting this relationship.

Our results predict that the largest excursions occur where the recovery time of Ω is shorter. As an indication of the degree of undersaturation, a site with more rapid recovery should be a site with less of a decrease in CaCO₃ wt% or, in terms of Fig. 3-10, lower $\Delta_n CaCO_3$. On the other hand, results predict larger excursions in the absence of bioturbation. Because bioturbation results in higher CaCO₃ wt% for a given pulse size than when bioturbation is not occurring, this implies that larger excursions should occur with higher $\Delta_n CaCO_3$. Therefore, the relationship between CaCO₃ content changes and the $\delta^{13}C$ excursion magnitude is unlikely to be a simple one for a suite of sites in which some are bioturbated, and some are not, or cease to be during the PETM event.

In this light, one way to view Figs. 3-1 and 3-10 is that they show two response regimes to the PETM carbon pulse: On the right hand side are sites for which the degree of perturbation of Ω is the main control on recovery time, and therefore on excursion magnitude. On the left are sites for which the rate of bioturbational mixing is the more important control on recovery time and excursion magnitude. Observations of lithologies of sites in Fig. 3-1 tend to support this scheme. On the right hand side, the sites for which we predict that recovery from undersaturation is the most important control include Walvis Ridge, for which bioturbation is reported to be nearly absent during the PETM at Site 1266 (Shipboard Scientific Party, 2004c), and absent for Sites 1262 and 1263

(Shipboard Scientific Party, 2004a, b). Diminished bioturbation is reported at Site 1001 for the PETM (Bralower et al., 1997). Laminae are present at Site 1220 during the PETM (Shipboard Scientific Party, 2002a). At Site 1221, bioturbation is present as burrows, and diminishes downcore within the PETM section (Shipboard Scientific Party, 2002b). In contrast, on the left hand side of Fig. 3-1, the Shatsky Rise sites and Site 690 display bioturbation throughout the PETM interval (Bralower et al., 1995; Thomas and Shackleton, 1996; Colosimo et al., 2006).

There is some support to be gained for this interpretation from the timescales of observed profiles of bulk carbonate δ^{13} C and CaCO₃ wt% from Walvis Ridge (Zachos et al., 2005). In Fig. 3-11, profiles are shown for Site 1262 (3600 m paleodepth), and Site 1263 (1500 m paleodepth) in blue. Model synthetic sediment cores (black lines) from depths of 3,600 m and 1,800 m in the Walvis Ridge region are superimposed. For Site 1262, the 2,300 Pg C scenario with bioturbation (Fig. 3-11A) comes the closest to matching the magnitude of the δ^{13} C excursion, but the onset and recovery of the excursion are much faster for the model sediment core than in the observed profile. The 21,000 and 46,000 Pg C cases come closer to matching the initial rate of the excursion, but do not produce an excursion of the appropriate magnitude (Fig. 3-11J,N). The 46,000 and 70,000 Pg C cases come closest to matching the CaCO₃ wt% profile at 1262 (Fig. 3-11P). For 1263, the shallower of the two sites, the 7,000 Pg C cases (Fig. 3-11G) come closest to matching the rate of the onset of the δ^{13} C excursion, but none of the model scenarios reproduces the excursion magnitude.

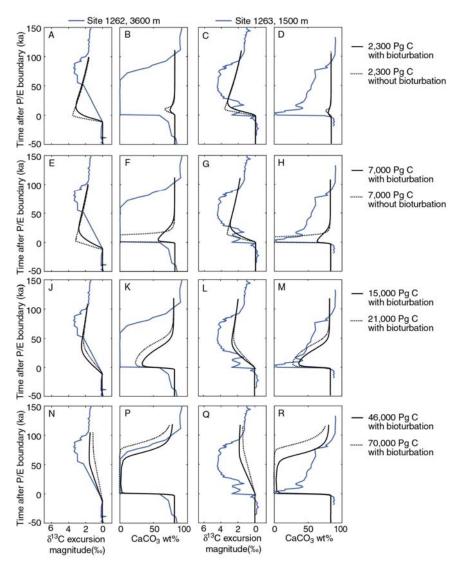


Figure 3-11: Comparison of sediment-core results with observations from Walvis Ridge by Zachos et al. (2005).

The 7,000 Pg C case without bioturbation (Fig. 3-11H) comes closest to reproducing the initial sharp decrease and recovery in the CaCO₃ wt% record, but no simulation reproduces the subsequent slower recovery. In summary, pulses of 21,000 Pg C and larger offer the better match to the CaCO₃ wt% record and to the rate of the δ^{13} C

excursion for the deeper site, 1262, but do not generate a large enough $\delta^{13}C$ excursion; for the shallower site, 1263, the 7,000 Pg C scenarios appear to offer a better match to the initial drop and recovery of the CaCO₃ wt% record and to the rate of onset of the $\delta^{13}C$ excursion, but do not provide a large enough $\delta^{13}C$ excursion. Despite the apparent lack of a consistent answer about the appropriate pulse size on the part of the model, it is interesting to note that the scenarios with timescales that come closest to matching the timescales of the Walvis Ridge records are also the scenarios that plot in the same locations as Sites 1262 and 1263 in Fig. 3-10.

Overall, the model-data comparison in Fig. 3-11 does present some conflicting results. First, the best match to the Site 1262 CaCO₃ wt% record comes from model scenarios with bioturbation, but evidence cited above indicates that there was little to no bioturbation at Walvis Ridge during the PETM. One explanation is that bioturbation decreased or stopped altogether at the PETM, but by the time the Ω recovered to allow burial of CaCO₃, bioturbation had recovered. This idea was offered by Ridgwell (submitted) to explain the apparent change in CaCO₃ wt% recovery rates at Site 1263 and others at Walvis Ridge.

A second conflicting result is the pulse sizes suggested by the sites at Walvis Ridge. At Site 1263, the model-data comparison suggests that the 7,000 Pg C pulse comes very close to simulating the PETM record. At Site 1262, however, it appears that a pulse closer to 70,000 Pg C than 7,000 Pg C would be a better approximation of what happened at the PETM. This may simply reflect the expected top-down recovery of the CCD, although a slow-down of ocean circulation could also contribute to the slower rate of recovery at the deeper site.

The model results shown here do not provide a clear-cut answer regarding the PETM., but they are informative nonetheless. The persistent underprediction of the $\delta^{13}C$ excursion in planktonic foraminifera by our model suggest that the excursion at the PETM was larger than -4‰, consistent with the larger $\delta^{13}C$ excursion in the terrestrial record. The difference between the marine and terrestrial $\delta^{13}C$ excursion magnitudes may be due to the fact that the entire $\delta^{13}C$ excursion cannot be recorded in marine $CaCO_3$ due to undersaturation, but that the terrestrial record has no similar limitations. The terrestrial $\delta^{13}C$ record would, however, be affected by the dissolution of old $CaCO_3$ because the higher $\delta^{13}C$ carbon would eventually enter the atmosphere. Therefore, mass balance calculations performed using the excursion recorded in the terrestrial record would likely underestimate either the $\delta^{13}C$ or the size of the PETM carbon pulse.

The behaviours of model sediments for different pulse sizes and with and without bioturbation suggest that while the changes in CaCO₃ wt% or δ^{13} C at a particular site may not be reliable indicators of pulse size, the timescale of the response might be. It remains to test the sensitivity of model timescales to weathering feedbacks and to different rates of mixing within the sediments, but in general, the similar range of timescales at Walvis Ridge and in our results is promising. If the timescale at Site 1263 at Walvis Ridge were a good indicator that the pulse at the PETM was ~7,000 Pg C, and if the terrestrial δ^{13} C excursion were representative of the PETM excursion, then given a system δ^{13} C of 0.5%, the pulse would have had δ^{13} C of -28 to -39%, suggesting a thermogenic methane component.

Conclusion

The magnitude of the PETM δ^{13} C excursion has been key to estimates of the size of the carbon pulse added to the global carbon-cycle at the Paleocene-Eocene boundary. The magnitude of the excursion is uncertain, however. While marine carbonate sediments record δ^{13} C excursions of -3 to-4% on average, δ^{13} C excursions recorded on land tend to be -5 to -7‰. There are several processes which occur after low δ^{13} C CO₂ mixes with surface ocean DIC and before CaCO₃ is preserved in the sediments that affect how well the original signal from DIC is propagated and preserved in the rock record. Our model results suggest that a particularly important consideration is the timescale of recovery of the saturation state of seawater with respect to calcite, where larger carbon pulses result in longer timescales. While seawater remains undersaturated, no net deposition of $CaCO_3$ occurs to record the $\delta^{13}C$ excursion in surface ocean DIC. The longer this condition persists, the further the δ^{13} C of DIC recovers toward pre-excursion values, and the smaller the excursion that will be recorded when CaCO₃ deposition resumes. Bioturbation is also a key factor, particularly when deposition resumes upon a low-CaCO₃ dissolution surface. By mixing older, higher δ^{13} C CaCO₃ with low δ^{13} C CaCO₃ deposited during the PETM, bioturbation results in a record that shows higher δ¹³C values and higher CaCO₃ content than actually existed. The absence of similar "pause" and "rewind" mechanisms on land may explain the difference between terrestrial and marine PETM records.

We conclude, therefore, that the marine $\delta^{13}C$ excursion and the decrease in marine CaCO₃ wt% are not consistently reliable indicators of the amount of low $\delta^{13}C$ carbon

added to the Earth system to generate the PETM. The timescales of the system recovery as indicated by δ^{13} C and CaCO₃ wt% recovery in marine sediments might be informative, however. The shortest timescales could be used to place an upper bound on pulse size. The seemingly conflicting views of the size of the PETM pulse from the comparison of model sediment cores with δ^{13} C and CaCO₃ wt% data from Walvis Ridge likely stems in part from the expected top-down recovery of the CCD, which is a function of circulation patterns and weathering rates.

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 $\label{eq:Appendix A} Appendix \, A$ Paleodepth and CaCO3 content data used in this study

Data plotted in Figs. 2-1B and 2-3 are listed under the appropriate sub-headings.

Site	<u>Paleodepth</u>			Sediment CaCO ₃ wt%		
	Depth (m)	Source	Baseline	PETM	Source	
Pacific Ocea	n					
164	5220	Van Andel (1975)	hiatus*	_	Van Andel (1975)	
166	4880	Van Andel (1975)	hiatus*	_	Van Andel (1975)	
170	5250	Van Andel (1975)	hiatus*	_	Van Andel (1975)	
303	5540	Van Andel (1975)	hiatus*	_	Van Andel (1975)	
304	5600	Van Andel (1975)	hiatus*	_	Van Andel (1975)	
307	5520	Van Andel (1975)	hiatus*	_	Van Andel (1975)	
313	3200	Van Andel (1975)	present	_	Van Andel (1975)	
316	4020	Van Andel (1975)	hiatus*	_	Van Andel (1975)	
1217	3200	Rea and Lyle (2005)	present	_	SSP (2002f)§	
LL4-GPC3	_		absent	_	Kyte et al. (1993)	
Sha	tsky Rise (a)					
577	1900	Miller et al. (1987)	96	hiatus	Bralower et al. (2002) Quillévéré et al. (2002)	
1209	2000 - 2500	SSP (2002a) [§]	96	84	Colosimo et al. (2006)	
1210	2000 - 3000	SSP (2002b)§	92	86	Colosimo et al. (2006) Colosimo et al. (2006)	
1211	2000 - 3000	SSP (2002c)§	95	78		
1212	2000 - 3000	SSP (2002d)§	98	83	Colosimo et al. (2006)	
Cen	tral Equatorial	Pacific (b)				
163	4490	Van Andel (1975)	absent	_	Van Andel (1975)	
165	4570	Van Andel (1975)	absent	_	Van Andel (1975)	
315	4120	Van Andel (1975)	hiatus*	_	Van Andel (1975)	
865	1300 – 1500	Thomas and Shackleton (1995)	96	93	Thomas et al. (1999) SSP (1995) [§]	
1220	2900	Rea and Lyle (2005)	90	0	SSP (2002g)§	
1221	3200	Rea and Lyle (2005)	74	3	Murphy et al. (2006)	
168, 169	5000	Van Andel (1975)	hiatus*	_	Van Andel (1975)	
Eas	tern Equatorial	Pacific (c)				
1215	3200	Rea and Lyle (2005)	74	-	SSP (2002e)§	
Caribbean S	Sea					
999	1750	Thomas et al. (1999)	61	0	Bralower et al. (1997)	
1001	1500	Thomas et al. (1999)	45	0	Bralower et al. (1997)	

This table is continues on the next page.

Table A: continued

Site		Sediment CaCO ₃ wt%				
	Depth (m)	Source	Baseline	PETM	Source	
Atlantic Oc	ean					
527	3400	Moore et al. (1984)	83	0	Thomas et al. (1999)	
549	2000 - 2300	Masson et al. (1985)	51	1	Thomas & Bralower (2005)	
1050	_	_	63	57	Rudnicki et al. (2001)	
1051	1000 - 2000	SSP(1998)§	56	52	Rudnicki et al. (2001)	
Wa	lvis Ridge (c)					
1262	3600	Zachos et al. (2005)	88	1	Zachos et al. (2005)	
1263	1500	Zachos et al. (2005)	88	1	Zachos et al. (2005)	
1266	2600	Zachos et al. (2005)	85	4	Zachos et al. (2005)	
1267	3200	Inferred [†]	80	1	Zachos et al. (2005)	
Southern O	cean, Atlantic s	ector (e)				
689	1400 – 1650	Kennett & Stott (1991)	91	_	SSP (1988)§	
690	2100	Zachos et al. (1993)	84	60	Thomas et al. (1999)	
Indian Ocea	an					
212	6050	Van Andel (1975)	hiatus*	_	Van Andel (1975)	
215	3450	Van Andel (1975)	present	_	Van Andel (1975)	
235	4300	Van Andel (1975)	absent	_	Van Andel (1975)	
236	2740	Van Andel (1975)	present	_	Van Andel (1975)	
240	3200	Van Andel (1975)	present	_	Van Andel (1975)	
241	3990	Van Andel (1975)	absent	_	Van Andel (1975)	
245	3560	Van Andel (1975)	present	_	Van Andel (1975)	
250	5080	Van Andel (1975)	hiatus*	_	Van Andel (1975)	
259	4430	Van Andel (1975)	present	_	Van Andel (1975)	
260	5080	Van Andel (1975)	hiatus*	_	Van Andel (1975)	
261	5400	Van Andel (1975)	hiatus*	_	Van Andel (1975)	
Ind	ian Ocean (f)					
211	4900	Van Andel (1975)	absent	_	Van Andel (1975)	
213	3520	Van Andel (1975)	present	_	Van Andel (1975)	
239	4220	Van Andel (1975)	absent	_	Van Andel (1975)	
248	4170	Van Andel (1975)	hiatus*	_	Van Andel (1975)	
256	5000	Van Andel (1975)	hiatus*	_	Van Andel (1975)	
738	1350	Barrera & Huber (1991)	90	70	SSP (1989)§	
766	3476	SSP (1990)§	93	_	SSP (1990) [§]	
1135	1000 - 2000	Inferred [‡]	95	_	SSP (2000) [§]	
1138	600 - 1200	Quilty (2002)**	93	_	SSP (2000)§	

^{*} Due to non-deposition

[§] SSP = Shipboard Scientific Party

[†] At present, Site 1262 is 400 m deeper than 1267. To estimate a paleodepth for 1267, 400 m was subtracted from the paleodepth of Site 1262.

[†] Inferred from paleodepths of surrounding sites, including the Paleocene paleodepths of Sites 738 (Barerra and Huber, 1991), 747, and 748 (Mackensen and Berggren, 1992), and Late Maastrichtian paleodepths of Sites 747, 748, 750 (Quilty, 1992), and 1138 (Quilty, 2002), and given the bathymetric divisions of Berggren and Miller (1989).

^{**} Late Campanian to Late Maastrichtian paleodepth.

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Appendix B

Model configuration

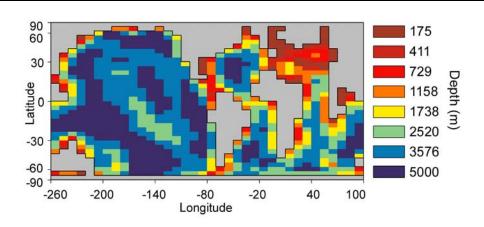


Figure B-1: Early Eocene bathymetry

Appendix C

Evaluation of goodness-of-fit between model output and observations

We address uncertainty regarding the Late Paleocene global weathering rate and the ratio of CaCO₃ to particulate organic carbon in the export of biogenic detritus from the surface ocean (the CaCO₃:POC rain ratio) by performing a 2-parameter sensitivity analysis comprising an ensemble of simulations with weathering fluxes of 30, 50, and 70 Tmol a⁻¹ of HCO₃ and spatiotemporally uniform rain ratios of 0.100, 0.175, and 0.250 Each ensemble member was spun up for 150 ka to reach geochemical equilibrium between weathering and sedimentation. The model CaCO₃ wt% output for each run is plotted in map view in Fig. C-1 with observations (sources listed in Appendix A) superimposed.

The best fit between model output and observed CaCO₃ wt% was identified by determining the lowest root mean squared error (RMSE) between model output and observations,

$$RMSE = \sqrt{\frac{1}{n}\Sigma(result - observation)^2}$$
 Eq. C-1

where n is the number of comparisons. First, CaCO₃ wt% was plotted versus depth for each of the model outcomes and the observations for three relatively data-rich regions: the central equatorial Pacific, Walvis Ridge in the Atlantic Ocean, and the southern Indian Ocean (identified as a, b, and c in Fig. 2-1A in the text). The profiles were fit using a linear weighted regression scheme described by Cleveland and Devlin (1988),

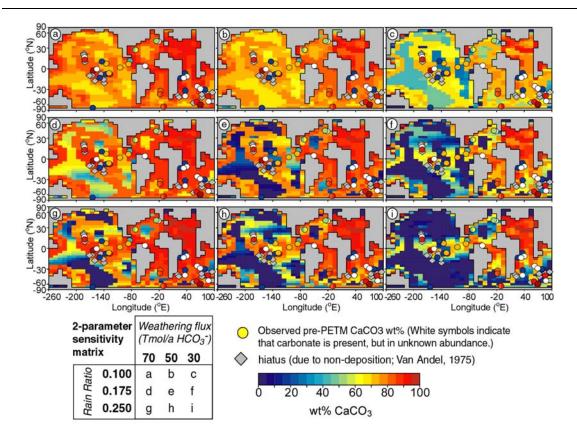


Figure C-1: Coretop CaCO₃ wt% output (maps) for an ensemble of model runs (panel labels correspond to the table in the legend) with observations superimposed (circle and diamond markers).

and available on-line in a data visualization toolbox (version 1.1) from Datatool (www.datatool.com) as the script loess.m for Matlab. This script requires that the user specify a set of points at which to compute the value of the fitted function, here set at 200 m increments; a smoothing parameter, α ; and the order of the polynomial to be used by the Matlab function polyfit, λ . Values of α and λ are listed in Table C-1.

Next, the calculated fits to both model output and observations were evaluated at model layer depths encompassed by the paleodepths of observations. The RMSE was calculated for each region (Fig. C-2 A,B,C) and the total RMSE was computed (Fig. C-2D).

Table C-1: Parameters for loess.m

Case to fit	Smoothing (α)	Polynomial order (λ)
Observations		
central equatorial Pacific Ocean		
Initial	0.5	2
PETM	0.65	2
Walvis Ridge, southern Atlantic Ocean		
Initial	1	1
PETM	1	1
southern Indian Ocean		
Initial	0.75	1
Model results		
Initial conditions		
50 Tmol HCO ₃ a ⁻¹ , CaCO ₃ :POC of 0.175 and 0.250	0.75	1
in the southern Indian Ocean		
70 Tmol HCO_3^- a ⁻¹ , CaCO ₃ :POC of 0.175 in the	1	2
southern Indian Ocean		
All other ensemble members and basins	1	1
PETM conditions	1	1

The best overall model-data fit (i.e., the scenario with the lowest total RMSE) is obtained with a $CaCO_3$:POC rain ratio of 0.175, higher than the mean value of \sim 0.14 characterizing the modern ocean (Ridgwell et al., 2007), and a total global weathering rate of 50 Tmol/a of HCO_3^- , which is larger than the 23-39 Tmol/a of HCO_3^- estimated for the modern (Munhoven, 2002) but less than that from early Eocene model simulations (Gibbs et al., 1999). The best overall model-data fit has a global average $CaCO_3$ wt% of 61%.

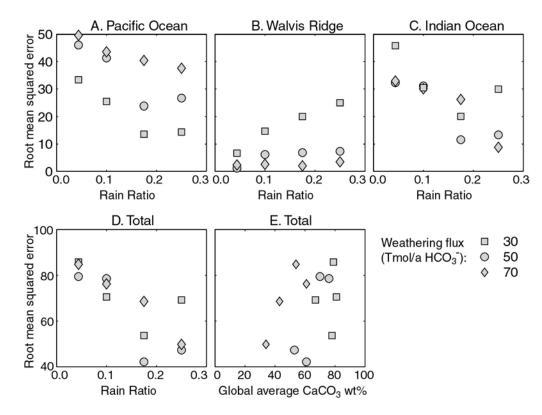


Figure C-2: Root mean squared errors (RMSE) for the 2-parameter sensitivity ensemble. The lowest overall RMSE (D, the sum of RMSE for data rich regions in A, B, and C) is achieved with a weathering flux of 50 Tmol/a of HCO₃⁻ and a rain ratio of 0.175. This scenario corresponds to a global average model CaCO₃ wt% of 61% (E).

The same analysis was performed to evaluate the best fit to observed PETM CaCO₃ wt% distributions (Fig. 2-6), except that the southern Indian Ocean region was not included because there are insufficient observations to generate a meaningful fit.

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